

Problem Set 6

Paramagnetism of localized magnetic moments

1 Classical treatment

In this first exercise we aim at calculating the contribution of localized magnetic moments to the magnetic susceptibility χ . We consider that the magnetic moments are fixed at the lattice sites of the crystal and take them as distinguishable and independent (*i.e.*, one given magnetic moment does not interact with the others).

- (a) To check the validity of the assumption of independence of magnetic moments, one can compare the potential energy (*i.e.*, the Zeeman energy)

$$E_Z = -\boldsymbol{\mu} \cdot \mathbf{B} \\ \sim \mu_B B$$

of a magnetic moment $\boldsymbol{\mu}$ (which is of the order of the Bohr magneton μ_B) in a magnetic field $\mathbf{B} = B \hat{z}$ to the dipole interaction energy. The latter is given by

$$E_{\text{dip}} = \frac{\mu_0}{4\pi r_{ij}^3} [\boldsymbol{\mu}_i \cdot \boldsymbol{\mu}_j - 3(\boldsymbol{\mu}_i \cdot \hat{r}_{ij})(\boldsymbol{\mu}_j \cdot \hat{r}_{ij})] \\ \sim \frac{\mu_0 \mu_B^2}{4\pi r_{ij}^3}$$

where $\boldsymbol{\mu}_i$ and $\boldsymbol{\mu}_j$ correspond to two magnetic moments located at the lattice sites \mathbf{r}_i and \mathbf{r}_j , respectively, with $\mathbf{r}_{ij} = \mathbf{r}_i - \mathbf{r}_j$.

Estimate the order of magnitude of these two energy scales (in eV) when r_{ij} is of the order of the typical distance between atoms in a solid and for a magnetic field $B = 1$ T. Compare also the Zeeman energy to the energy of thermal fluctuations at room temperature.

Values: $\mu_0 = 4\pi \times 10^{-7}$ Tm/A, $\mu_B = 9.3 \times 10^{-24}$ J/T, $1 \text{ eV} = 1.6 \times 10^{-19}$ J.

- (b) We now treat each (independent) magnetic moment classically and allow for arbitrary orientation of $\boldsymbol{\mu}$ with respect to the magnetic field \mathbf{B} . Express the Zeeman energy E_Z in terms of the absolute values of these vectors and the angle θ between them.
- (c) The probability density for a given orientation of the magnetic moment $\boldsymbol{\mu}$ in an external magnetic field is determined by its energy through

$$P(\boldsymbol{\mu}) = \frac{1}{Z} e^{-\beta E_Z(\theta)},$$

with $\beta = 1/k_B T$ the inverse temperature, and where Z is the (canonical) partition function given by the integral over all orientations

$$Z = \int d\Omega e^{-\beta E_Z(\theta)}.$$

Calculate Z .

- (d) The mean magnetization $\langle \mathbf{M} \rangle$ of a material that contains a density $n = N/V$ of N magnetic moments (with V the sample volume), under the influence of an external magnetic field, is defined as the total average magnetic moment per unit volume. Show that $\langle \mathbf{M} \rangle = n \langle \boldsymbol{\mu} \rangle$. Calculate explicitly the average magnetization and show that $\langle \mathbf{M} \rangle = \langle M_z \rangle \hat{z}$ with $\langle M_z \rangle = n \mu \mathcal{L}(\beta \mu B)$, where

$$\mathcal{L}(x) = \coth x - \frac{1}{x} \tag{1.1}$$

is the Langevin function.

- (e) Discuss the dependence of the magnetization in the limits of low field/high temperature $\mu B/k_B T \ll 1$ and strong field/low temperature $\mu B/k_B T \gg 1$. Sketch $\langle M_z \rangle$ as a function of $\mu B/k_B T$.
- (f) The zero-field magnetic susceptibility is defined as

$$\chi = \left. \frac{\partial \langle M_z \rangle}{\partial H} \right|_{H=0},$$

where $\mathbf{H} = \mathbf{B}/\mu_0 - \mathbf{M}$. Argue that

$$\chi \simeq \mu_0 \left. \frac{\partial \langle M_z \rangle}{\partial B} \right|_{B=0}.$$

- (g) Calculate the magnetic susceptibility using the low-field limit from Question (e) and show that it is always paramagnetic and follows the Curie law $\chi = C/T$, where C is a constant to be specified.

2 Quantum treatment

In this second exercise, we still consider a system of N independent atoms in a volume V , each of which has a total angular momentum (in units of \hbar) $\mathbf{J} = \mathbf{S} + \mathbf{L}$, with \mathbf{S} and \mathbf{L} the spin and orbital angular momentum, respectively. The absolute square value $\mathbf{J}^2 = J(J+1)$, where J is the quantum number of the total angular momentum. The magnetic moment of an atom is $\boldsymbol{\mu} = -g\mu_B \mathbf{J}$, where g is the Landé factor of the considered atoms.

- (a) Write the Hamiltonian describing the Zeeman energy of an atom in an external magnetic field \mathbf{B} parallel to the z axis. Specify the possible values J_z of the projection of the angular momentum \mathbf{J} of an atom on the z axis.
- (b) Write the (canonical) partition function Z of each atom. Show that it can be expressed as the ratio of two hyperbolic sines.¹
- (c) The free energy (per atom) is defined as $F = -k_B T \ln Z$. Show that the mean value of the magnetic moment is given by $\langle \mu_z \rangle = -\partial F / \partial B$.
- (d) Show that the average magnetization in z direction $\langle M_z \rangle$ defined as the mean total magnetic moment per unit volume can be expressed as

$$\langle M_z \rangle = M_s B_J(\beta g \mu_B J B), \quad (2.1)$$

where

$$B_J(x) = \frac{2J+1}{2J} \coth\left(\frac{2J+1}{2J}x\right) - \frac{1}{2J} \coth\left(\frac{1}{2J}x\right) \quad (2.2)$$

denotes the Brillouin function. What is the expression of the saturation magnetization M_s ?

- (e) Let us first consider the case of vanishing orbital moment and spin $1/2$, *i.e.*, $J = 1/2$. Show that $B_{1/2}(x) = \tanh x$ and discuss then the result of Eq. (2.1).
- (f) We now consider the classical limit, in which $\hbar \rightarrow 0$, which amounts to consider $J \rightarrow \infty$. Demonstrate that $B_\infty(x) = \mathcal{L}(x)$, where $\mathcal{L}(x)$ is the Langevin function from Eq. (1.1).
- (g) We finally consider the case of an arbitrary, finite J . Sketch the Brillouin function (2.2) as a function of the parameter x for various values of J .
- (h) Still for finite J , deduce the value of the zero-field susceptibility χ . Is the system paramagnetic?

¹Note that

$$\sum_{k=0}^n q^k = \frac{1-q^{n+1}}{1-q}, \quad (q \neq 1).$$